



A 700 pc Extended Coronal Gas Emission in the Circinus Galaxy

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Abstract

We report the first characterization of an extended outflow of high ionized gas in the Circinus Galaxy by means of the coronal line (CL) [Fe VII] $\lambda 6087$. This emission is located within the ionization cone already detected in the [O III] $\lambda 5007$ line and is found to extend up to a distance of ~ 700 pc from the active galactic nucleus. The gas distribution appears clumpy, with several knots of emission. Its kinematics is complex, with split profiles and line centroids shifted from the systemic velocity. The physical conditions of the gas show that the extended coronal emission is likely the remnant of shells inflated by the passage of a radio jet. This scenario is supported by extended X-ray emission, which is spatially coincident with the morphology and extension of the [Fe VII] $\lambda 6087$ gas in the NW side of the galaxy. The extension of the coronal gas in the Circinus galaxy is unique among active galaxies and demonstrates the usefulness of CLs for tracing the shock ionization component in these objects.

Unified Astronomy Thesaurus concepts: Shocks (2086); Active galactic nuclei (16); Radio jets (1347); Line intensities (2084)

1. Introduction

The relevance of the kinetic channel as a major way of releasing nuclear energy to the ISM in low-luminosity active galactic nuclei (AGN) has been underestimated so far (Merloni & Heinz 2013; Wylezalek & Morganti 2018). The lack of convincing observational evidence of strong feedback in such a class of objects has been, up to very recently, scarce.

Recent works, however, reveal that low-power radio jets can play a major role in driving fast, multiphase, galaxy-scale outflows even in radio-quiet AGN (Rodríguez-Ardila et al. 2017; May et al. 2018; Jarvis et al. 2019). Mass outflow rates, of up to $8 M_{\odot} \text{ yr}^{-1}$, similar to that found in powerful radio-loud AGN have been derived. Overall, they all point toward the relevance of low Eddington sources in the delivery of kinetic energy to the ISM.

Traditionally, the identification of outflows in the warm, ionized phase of AGN has been done by means of the [O III] $\lambda 5007$ line (Greene et al. 2011). In combination with integral field spectroscopy, critical insights about the geometry, structure, extension, and physical conditions of the outflowing gas can be derived. This technique has allowed the discovery of kiloparsec-scale outflows in AGN samples (Humphrey et al. 2010; Harrison et al. 2014; Karouzos et al. 2016).

However, because [O III] is also emitted in the galaxy disk and star-forming regions, isolating the contribution due to outflows is not straightforward. In this respect, Rodríguez-Ardila et al. (2006) showed that high-ionization lines such as [Fe VII] $\lambda 6087$ in the optical or [Si VI] $1.963 \mu\text{m}$ in the near-infrared are excellent tracers of the ionized component of the outflows. Both lines are usually the brightest coronal lines (CLs) in their respective wavelength regimes. The energy required for their production ($\chi \geq 100$ eV, where χ is the ionization potential required to produce the ion) rules out stellar or galactic origins. The case of ESO 428-G 14 (May et al. 2018) is emblematic. These authors employed the [Si VI] to trace the jet-driven mechanical energy and the corresponding mass outflow deposited by the jet in the central 170 parsecs of that object.

At the adopted distance of 4.2 Mpc ($1'' \sim 20.4$ pc), Circinus is the closest Seyfert 2 galaxy to us. Because of its proximity, angular resolution on scales of a few tens of parsecs can be reached even at seeing-limited conditions. This makes Circinus an excellent laboratory to study the physics of the ionized gas in a large range of distances. This includes the close examination of one of the most prominent characteristics of this object: a spectacular one-sided ionization cone, seen mostly in the NW side of the galaxy using the [O III] $\lambda 5007$ line that extends for at least $30''$ (Elmouttie et al. 1998).

In addition, Circinus is widely known for its prominent CL spectrum (Moorwood et al. 1996). Prieto et al. (2005) determined the size and morphology of the coronal line region (CLR) for the first time in Circinus by means of adaptive (AO) imaging observations. Using [Si VII] $2.48 \mu\text{m}$ as a tracer of the CLR, they found that the region emitting that line extends from the nucleus up to 30 pc. Later, Müller-Sánchez et al. (2006, 2011) using *K*-band SINFONI/Very Large Telescope (VLT) AO integral field unit (IFU) spectroscopy revealed extremely strong coronal emission lines of [Si VI] and [Ca VIII] as well as [Al IX]. In all three CLs, in addition to an unresolved, narrow component, they found a broad (FWHM $> 300 \text{ km s}^{-1}$) blueshifted component, which is spatially extended. It was argued that the narrow component arises in clouds physically close to the AGN while the blue wing originates from cloudlets that have been eroded from the main clouds and are accelerated outward. Due to the small field of view (FoV) of the detector ($3'' \times 3''$), it was not possible to determine if extended CL emission out of the nuclear region was present in Circinus.

Here, we report optical observations of the Circinus Galaxy focusing on two main aspects: (i) the determination of the full extension of the coronal gas and its relationship with galactic feedback; (ii) the role of a jet in shaping the morphology of the high-ionization gas in this object.

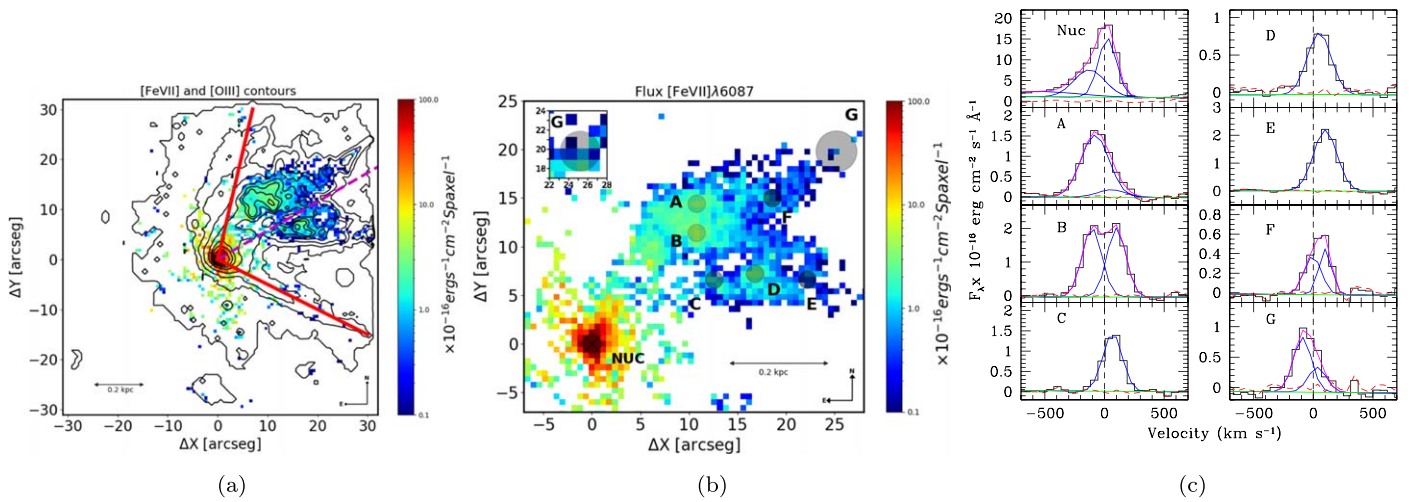


Figure 1. (a) Emission map of the [Fe VII] $\lambda 6087$, overlaid on [O III] $\lambda 5007$ contours. The red lines mark the edges of the ionization cone (Mingozi et al. 2019). The dashed magenta line indicates the PA = $295^\circ \pm 5^\circ$ of the radio continuum reported by Elmouttie et al. (1998). (b) Zoom-in on the NW ionization cone, emphasizing the extended [Fe VII] emission. The letters A to G indicate regions where the line profiles shown in panel (c) were extracted. The AGN position is labeled “Nuc.” The inset in the upper left corner shows region G, with the spaxels binned to a size of $1''.2 \times 1''.2$ to increase the S/N to >3 . (c) Selected [Fe VII] $\lambda 6087$ line profiles, in velocity space, detected in the extended coronal gas. The dashed vertical line marks the systemic velocity of the galaxy. In each panel, the observed profile (black histogram) was fit with Gaussian functions (blue lines). The total fit is in magenta, the continuum level is in green and the residual after subtraction of the total fit to the observed data is in dashed red.

2. Observations and Data Analysis

Data for Circinus, observed using MUSE/VLT was retrieved from the ESO science portal. They were obtained under an equivalent airmass of 1.3, average seeing of $0''.78$, and exposure time of 1844 s. The MUSE IFU cube consists of $\sim 300 \times 300$ spaxels, for a total of over 90,000 spectra with a spatial sampling of $0''.2 \times 0''.2/\text{spaxel}$. The FoV of $1' \times 1'$ covers the central part of the galaxy, spanning an area of $\sim 1.2 \times 1.2 \text{ kpc}^2$. The spectral resolution varies from 1750 at 4650 Å to 3750 at 9300 Å. The IFU cube, as retrieved from the science portal, is fully reduced, including calibration in flux (in absolute units) and wavelength.

The MUSE data cube was analyzed making use of a set of custom python scripts developed by us as well as software publicly available in the literature. First, we rebinned the cube to a spatial sampling of $0''.6 \times 0''.6$, reducing the total number of spaxels to $\sim 10,000$. This procedure was employed to increase the signal-to-noise ratio (S/N) of weak emission lines within the ionization cone. As the angular extension of the features we plan to study are of several arcseconds in size, the new spaxel scale does not affect the results reported here.

We then remove the stellar continuum across the whole spectral range of MUSE (4700–9100 Å) using STARLIGHT (Cid-Fernandes et al. 2005) and the Bruzual & Charlot (2003) stellar libraries.

After this procedure, we measured the integrated emission line fluxes of $H\alpha$ and $H\beta$ at every spaxel to determine the extinction (Galactic and intrinsic) affecting the gas. This was done assuming an intrinsic line ratio $H\alpha/H\beta = 3.1$ and the Cardelli et al. (1989) extinction law. Our results point out a Galactic extinction, A_V , of 2.1 mag, in agreement with that reported by Mingozi et al. (2019). The internal extinction varies from nearly 0 to 2 mag. In the inner region of the ionization cone, A_V is rather homogeneous and low, with values between 0.1 and 0.3 mag. Integrated line fluxes measured at each spaxel were corrected by the extinction measured accordingly.

Finally, for each spaxel, we fit the [Fe VII] $\lambda 6087$ line with a single or a combination of Gaussian functions. Typically, one or two components were necessary to reproduce the observed profile. This procedure allowed us to measure the flux, velocity, and velocity dispersion of the line across the whole MUSE FoV.

3. Morphology and Kinematics of the Coronal Gas

A careful inspection of each spaxel in the data cube of the Circinus Galaxy reveals a very rich emission line spectrum, with prominent [O III] $\lambda 5007$ and $H\alpha$ lines across most of the FoV covered by MUSE. Mingozi et al. (2019) recently presented the gas distribution for these lines using the same data set as here. For instance, the $H\alpha$ disk component flux map can be seen overlaid with the [O III] $\lambda 5007$ outflow component as well as the ISM properties in the Circinus galaxy (see their Figures 2 and 3, respectively). In this Letter, we concentrate on the 2D properties of the extended coronal gas of that AGN. Oliva et al. (1999), using optical spectroscopy, had already reported extended [Fe VII] emission up to $22''$ from the center at a PA = 318° . Mingozi et al. (2019) reported the presence of extended [Fe VII] $\lambda 6087$ associated to the ionization cone, but no information about its full extension, morphology, and gas physical properties were presented.

Figure 1 shows, for the first time in the literature, the most complete [Fe VII] $\lambda 6087$ emission line map from the MUSE data cube. White areas correspond to masked regions with S/N < 3 . It can be seen that the central region of the NW ionization cone is filled with a conspicuous, extended coronal emission. The size of the highly ionized region occupies at least a projected area of $400 \times 300 \text{ pc}^2$. The gas distribution appears clumpy, arranged in several condensations, most of them contained within two main regions of emission, the latter two separated by a strip of very little or no [Fe VII] and coincident with the PA of the radio continuum ($295^\circ \pm 5^\circ$, Elmouttie et al. 1998). In the following, we will refer to this radio continuum emission as the radio jet. The region to the north (hereafter NR), is elongated in the SE–NW direction,

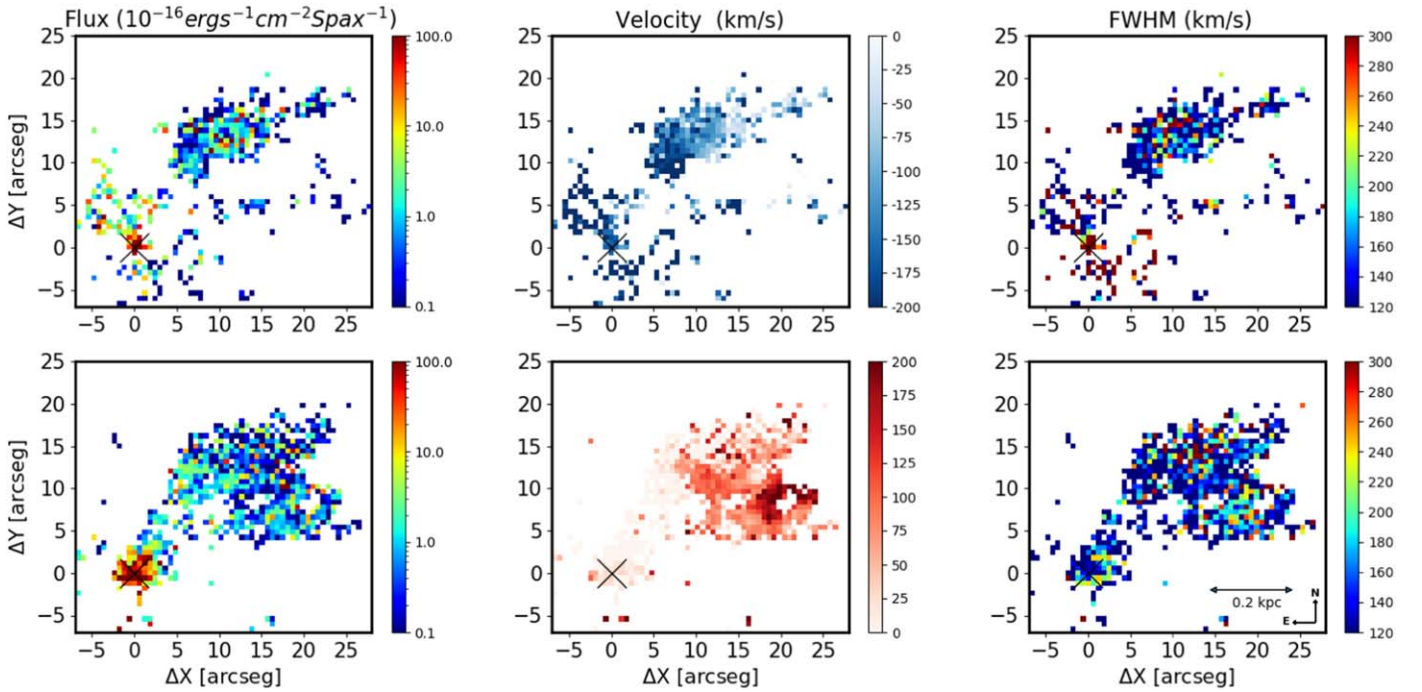


Figure 2. Flux distribution (left panel), velocity map (center panel), and FWHM of the line (right panel) for the blue and red components (upper and lower panels, respectively) of the [Fe VII] 6087 Å line. The cross marks the position of the AGN.

measuring approximately 200 pc SE–NW and 135 pc NE–SW. It contains a bright spot (which Oliva et al. 1999 named “knot C”), at $\sim 2''$ off-center of NR, toward the bottom edge of the cloud. Here, it is identified by the letter B (see the central panel of Figure 1).

The region to the south (hereafter SR) is composed of three smaller blobs of emission, (identified by the letters C to E in the central panel of Figure 1) arranged nearly in the E–W direction.

In addition to the NR and SR, clouds of smaller size emitting [Fe VII] (< 20 pc in radius) are also visible in the image. Most of them are located outwards of the far edge of the NR. The farthest one is detected up to $\sim 34''$ from the nucleus (~ 700 pc from the AGN). This spot is labeled G in the central panel of Figure 1. In order to clearly detect [Fe VII] at an S/N > 3 in G, we further re-binned the spaxels to $1''2 \times 1''2$ in size and integrated the flux within a circular aperture of $4''$ in radius. In the remainder of this manuscript, we will name the region enclosed by points A to G in Figure 1 as the extended emission. We do recognize the presence of extended, circumnuclear emission within a radius of $\sim 5''$ centered at the nucleus, but that is not the main focus of this work.

From the MUSE data, the integrated [Fe VII] flux in the circumnuclear region is $49.4 \times 10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2}$. The extended CL emission, integrated across the whole NW region (see panel (b), Figure 1), amounts to $6.5 \times 10^{-14} \text{ erg s}^{-2} \text{ cm}^{-2}$. These values are already corrected for extinction. Thus, the extended emission represents more than 10% of the nuclear value.

The [Fe VII] 6087 Å emission line profiles display a complex structure, with splitted lines at most spaxels in the NR. The relative intensity between the peaks and the width of the individual components vary appreciably. This is shown in panel (c) of Figure 1, where [Fe VII] emission line profiles extracted at different positions of the extended emission are plotted. Gaussian fitting carried out on these lines evidences the

necessity of two components at positions A, B, F, and G in order to adequately reproduce the observed profiles. The line profiles at positions C, D, and E are mostly redshifted, with just the red peak visible. Cospatial to the [Fe VII] emission, lines of H I, [O III], [O II], [N II], [O I], [S II], and [S III] are also observed in the corresponding integrated spectra (not shown here).

Figure 2 shows the integrated flux (left panel), velocity (central panel), and FWHM (right panel) for the blue component (upper row) and red component (bottom row) of the [Fe VII] 6087 Å line. All FWHM values were previously corrected in quadrature for the instrumental resolution. The velocity pattern in the extended emission observed in Figure 2 rules out that the coronal gas is rotating with the galaxy disk. This result is in agreement with previous works on the gas kinematics within the ionization cone of Circinus, pointing out that most of the material in that structure is out of the galaxy plane (Elmouttie et al. 1998).

Figure 2 also shows that the FWHM of both blue and red components are in the interval $\sim 160\text{--}200 \text{ km s}^{-1}$. At some spaxels, though, it jumps to $\sim 260 \text{ km s}^{-1}$ or higher values. At 80% of the line width, the velocity reaches $\sim 350 \text{ km s}^{-1}$ (see panel (c) in Figure 1). Moreover, the relative separation between the peaks varies along the extended emission, from 130 km s^{-1} at regions F and G to nearly 250 km s^{-1} at A and B.

It is also interesting to see that the bulk of the coronal gas shows evidence of expanding shells. For the blue component, the shift of line decreases from $\simeq -200 \text{ km s}^{-1}$ in the region that is closest to the central source, to the systemic velocity at the outermost portion of the cloud. Thus, if the gas is traveling toward the observer, we see preferentially the portion closest to us. For the red peak, receding velocities of up to $\sim 200 \text{ km s}^{-1}$ are detected. These clouds are located preferentially in the regions along the PA of the radio jet.

4. Discussion

Most previous detections of extended coronal gas in low-luminosity radio-quiet AGN resolved CL emission at scales of <300 pc (Müller-Sánchez et al. 2011; Rodríguez-Ardila et al. 2017; May et al. 2018). Ramos Almeida et al. (2017) reported that the CLR in the teacup galaxy as measured from the narrow component of the line is unresolved (<760 pc), but the outflowing component is resolved and its seeing-deconvolved radial size is 860 pc. In this radio-quiet quasar, there seems to be evidence for a jet driving the ionized outflow. Extended CL gas at kiloparsec scales has already been observed in radio galaxies. Tadhunter et al. (1988) presented solid evidence of [Fe VII] $\lambda 6087$ emission from a cloud at ~ 8 kpc from the AGN in PKS 2152-69. In that source, as well as in other similar radio-loud AGN (Solórzano-Iñarrea & Tadhunter 2003), the effect of induced shocks by a jet/cloud interaction has been claimed as the origin of the extended coronal emission.

The coronal gas kinematics shown in the preceding section suggest a scenario where the radio jet inflates ISM gas located along the region where it propagates. Shocks produced by this interaction result in the observed coronal emission. Because of the high-ionization state, the gas is mostly optically thin, allowing us to see both the approaching and receding components of the shells.

In order to test the above hypothesis, we first explore if photoionization by the AGN is able to explain the large-scale H I emission detected cospatially to [Fe VII]. To this purpose, we estimate the luminosity predicted for the $H\alpha$ line assuming a pure photoionization scenario. The bolometric luminosity of the AGN in Circinus is $10^{10} L_{\odot}$ (Oliva et al. 1999). The solid angle covered by the NR and SR region as measured from the MUSE data is, respectively, 0.137 sr and 0.024 sr. Both solid angles summed represent a fraction of 0.01 of all the surface area seen by the AGN. The integrated luminosity of the $H\beta$ line can be estimated by Equation (1) (Osterbrock 1989):

$$L(H\beta) = h\nu_{H\beta}\eta_e^2 f V \alpha_{H\beta}^{\text{eff}}, \quad (1)$$

where h is the Planck constant, $\nu_{H\beta}$ is the line frequency of the $H\beta$, η_e is the electronic density, V is the volume emitted in the region, f is the filling factor, and $\alpha_{H\beta}^{\text{eff}}$ is the recombination coefficient.

Taking into account a covering factor of 5% in Circinus, we estimate a $H\beta$ luminosity of 2.2×10^{39} erg s $^{-1}$ (or $5.9 \times 10^5 L_{\odot}$). Using a theoretical ratio $H\alpha/H\beta = 3.1$, it translates to a $H\alpha$ luminosity of 6.9×10^{39} erg s $^{-1}$ ($1.8 \times 10^6 L_{\odot}$). From the observed integrated $H\alpha$ flux in the ionization cone cospatial to the [Fe VII] emission, we derived a luminosity of 5.9×10^{39} erg s $^{-1}$ ($1.5 \times 10^6 L_{\odot}$). Thus, the AGN is just barely able to power the $H\alpha$ emission cospatial to the coronal gas. As $H\alpha$ extends from the nucleus up to the edge of the data cube, we conclude that the AGN alone is not sufficient to account for the extended low-ionized emission. This gives further support to the shock-driven coronal gas scenario.

Mingo et al. (2012) had already shown that the AGN in Circinus is able to weakly ionize gas up to ~ 700 pc away, but strong ionization would only be produced within the inner ~ 200 pc. They showed that the extended emission observed in the [O III] $\lambda 5007$ line in Circinus, which coincides with that of [Fe VII] (see Figure 1, left panel) is most likely caused by a jet-

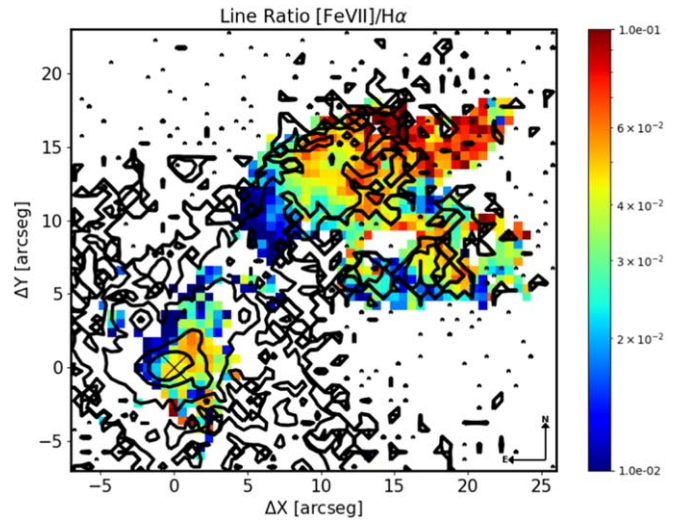


Figure 3. Extinction corrected emission line flux ratio [Fe VII] $\lambda 6087/H\alpha$ for the highest ionized portion of the cone. The region in white corresponds to values with $S/N < 3$ or where the [Fe VII] line is not detected. The cross marks the position of the AGN. The contours represent Chandra ACIS image in the 0.58 keV band from Smith & Wilson (2001).

driven outflow, which would be driving shells of strongly shocked gas into the halo of the host galaxy.

Additional support for this scenario can be obtained from the emission line flux ratio [Fe VII]/ $H\alpha$ (see Figure 3), which directly reflects the ionization state of the gas. Figure 3 shows that the values of that ratio increases outwards, varying from ~ 0.01 at the side of the NR facing the AGN (~ 77 pc) to ~ 0.1 in the parcel of gas located at 700 pc from the nucleus. A similar effect is also seen in the SR. Here, [Fe VII]/ $H\alpha$ increases by a factor of ~ 6 between the closest and farthest points from the AGN.

Oliva et al. (1999) invoked a large N/O overabundance to explain the rich emission line spectrum in “knot C” (cloud B, in our Figure 1). Although this is indeed a plausible explanation, notice that the [Fe VII]/ $H\alpha$ ratio measured in cloud B is nearly a factor of 3 smaller than that in Cloud G. Therefore, such a scenario fails at explaining the increasing gas excitation reported here. Assuming that the radio jet is launched from the nucleus, its motion across the ionization cone is highly supersonic, so high-velocity shocks are likely to contribute to the ionization of the line-emitting gas. Long after its passage through the ISM, the gas ionization should decrease due to cooling, while the gas most recently affected by the jet passage should display the largest excitation. The observed distribution of values of the [Fe VII]/ $H\alpha$ in Figure 3 is consistent with this scenario, as the gas excitation is the largest in the regions most distant from the AGN.

Theoretical models carried out by Mukherjee et al. (2018) and applied to IC 5063, a low-luminosity Seyfert 2 galaxy, shows a highly perturbed ISM by an expanding radio jet. As the jet floods through the inter-cloud channels of the disk, it ablates, accelerates, and disperses clouds from the central regions of the galaxy, in a similar manner as seen in Circinus. Moreover, models of Contini & Viegas (2001) show that shocks with velocities of 200 km s $^{-1}$ and preshock densities of $n_H = 200$ cm $^{-3}$ produce [Fe VII]/ $H\alpha$ ratios > 0.1 . Both the FWHM of the lines (see Figure 2) and the gas density in the region emitting [Fe VII], of a few hundred cm $^{-3}$ (see Mingozi et al. 2019) are consistent with this scenario.

Finally, the excellent agreement between the extended X-ray emission from the Chandra ACIS image in the 0.58 keV band (Smith & Wilson 2001) with the bulk of the [Fe VII] emission, shown in Figure 3, illustrates the strong interplay between AGN jet-driven shocks and the high-ionization lines, not seen before in an AGN. In effect, Mingo et al. (2012) interpreted this extended X-ray emission most likely caused by a jet-driven outflow, which is driving shells of strongly shocked gas into the halo of the host galaxy.

This work shows convincing evidence that the jet in the Circinus galaxy is inducing shocks capable of ionizing gas at large distances and possibly pushing outwards the ionized gas. The jet should drive the best resolved and most extended outflow of high-ionization gas ever observed in a low-luminosity AGN, extending up to 700 pc from the central engine. Our work also highlights the use of optical CL emission to gather kinematic information not available in the high-ionization gas observed in X-rays via Chandra ACIS data.

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References

Bruzual, G., & Charlot, S. 2003, *MNRAS*, 344, 1000

- Cardelli, J. A., Clayton, G. C., & Mathis, J. S. 1989, *ApJ*, 345, 245
- Cid-Fernandes, R., Mateus, A., Sodré, L., Stasińska, G., & Gomes, J. M. 2005, *MNRAS*, 358, 363
- Contini, M., & Viegas, S. M. 2001, *ApJS*, 132, 211
- Elmoultie, M., Koribalski, B., Gordon, S., et al. 1998, *MNRAS*, 297, 49
- Greene, J. E., Zakamska, N. L., Ho, L. C., & Barth, A. J. 2011, *ApJ*, 732, 9
- Harrison, C. M., Alexander, D. M., Mullaney, J. R., & Swinbank, A. M. 2014, *MNRAS*, 441, 3306
- Humphrey, A., Villar-Martín, M., Sánchez, S. F., et al. 2010, *MNRAS*, 408, L1
- Jarvis, M. E., Harrison, C. M., Thomson, A. P., et al. 2019, *MNRAS*, 485, 2710
- Karouzos, M., Woo, J.-H., & Bae, H.-J. 2016, *ApJ*, 819, 148
- May, D., Rodríguez-Ardila, A., Prieto, M. A., et al. 2018, *MNRAS*, 481, L105
- Merloni, A., & Heinz, S. 2013, in *Planets, Stars and Stellar Systems*, Vol. 6, ed. W. Keel (Dordrecht: Springer), 503
- Mingo, B., Hardcastle, M. J., Croston, J. H., et al. 2012, *ApJ*, 758, 95
- Mingozzi, M., Cresci, G., Venturi, G., et al. 2019, *A&A*, 622, 146
- Moorwood, A. F. M., Lutz, D., Oliva, E., et al. 1996, *A&A*, 315, L109
- Mukherjee, D., Wagner, A. Y., Bicknell, G. V., et al. 2018, *MNRAS*, 476, 80
- Müller-Sánchez, F., Davies, R. I., Eisenhauer, F., et al. 2006, *A&A*, 454, 481
- Müller-Sánchez, F., Prieto, M., Hicks, E., et al. 2011, *ApJ*, 739, 69
- Oliva, E., Marconi, A., & Moorwood, A. F. M. 1999, *A&A*, 342, 87
- Osterbrock, D. E. 1989, *Astrophysics of Gaseous Nebulae and Active Galactic* (Mill Valley, CA: Univ. Science Books)
- Prieto, M. A., Marco, O., & Gallimore, J. 2005, *MNRAS*, 346, L28
- Ramos Almeida, C., Piqueras López, J., Villar-Martín, M., & Bessiere, P. S. 2017, *MNRAS*, 470, 964
- Rodríguez-Ardila, A., Prieto, M. A., Mazzalay, X., et al. 2017, *MNRAS*, 470, 2845
- Rodríguez-Ardila, A., Prieto, M. A., Viegas, S., & Gruenwald, R. 2006, *ApJ*, 653, 1098
- Smith, D. A., & Wilson, A. S. 2001, *ApJ*, 557, 180
- Solórzano-Iñárrrea, C., & Tadhunter, C. N. 2003, *MNRAS*, 340, 705
- Tadhunter, C. N., Fosbury, R. A. E., Alighieri, S. S., et al. 1988, *MNRAS*, 235, 403
- Wylezalek, D., & Morganti, R. 2018, *NatAs*, 2, 181